

**Relationships involving D -functions and
Clebsch-Gordan Coefficients
—Version II—**

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abstract

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1 D -functions

For the rotation matrix, we use the definition as given in Rose [1], namely,

$$\begin{aligned} D_{m'm}^j(\alpha, \beta, \gamma) &= \langle jm' | e^{-i\alpha J_z} e^{-i\beta J_y} e^{-i\gamma J_z} | jm \rangle \\ &= e^{im'\alpha} d_{m'm}^j(\beta) e^{-im\gamma} \end{aligned} \quad (1)$$

By definition, the matrices $D_{m'm}^j$ are unitary and satisfy the group property:

$$\sum_k D_{m'k}^j(R) D_{mk}^{j*}(R) = \sum_k D_{km'}^j(R) D_{km}^{j*}(R) = \delta_{mm'} \quad (2)$$

$$D_{m'm}^j(R_2 R_1) = \sum_k D_{m'k}^j(R_2) D_{km}^j(R_1) . \quad (3)$$

The D -functions are normalized according to

$$\int dR D_{\mu_1 m_1}^{j_1*}(R) D_{\mu_2 m_2}^{j_2}(R) = \frac{8\pi^2}{2j_1 + 1} \delta_{j_1 j_2} \delta_{\mu_1 \mu_2} \delta_{m_1 m_2} , \quad (4)$$

where $R = R(\alpha, \beta, \gamma)$ and $dR = d\alpha d\cos\beta d\gamma$.

The functions $d_{m'm}^j$ have the following symmetry properties:

$$d_{m'm}^j(\beta) = (-)^{m'-m} d_{mm'}^j(\beta) \quad (5)$$

$$d_{m'm}^j(\beta) = (-)^{m'-m} d_{-m'-m}^j(\beta) \quad (6)$$

$$d_{m'm}^j(\pi - \beta) = (-)^{j+m'} d_{m'-m}^j(\beta) \quad (7)$$

$$d_{m'm}^j(\pi) = (-)^{j-m} \delta_{m', -m} \quad (8)$$

Owing to Eq. (6), the D -functions have the symmetry

$$D_{m'm}^{j*}(\alpha, \beta, \gamma) = (-)^{m'-m} D_{-m'-m}^j(\alpha\beta\gamma) . \quad (9)$$

One may use the identity

$$R(\pi + \alpha, \pi - \beta, \pi - \gamma) = R(\alpha, \beta, \gamma) R(0, \pi, 0) \quad (10)$$

to show that

$$D_{m'm}^j(\pi + \alpha, \pi - \beta, \pi - \gamma) = (-)^{j-m} D_{m'-m}^j(\alpha, \beta, \gamma) \quad (11)$$

and

$$D_{m'm}^j(\pi + \phi, \pi - \theta, 0) = e^{i\pi j} D_{m'-m}^j(\phi, \theta, 0) \quad (12)$$

or taking $\phi \rightarrow -\phi$ and using Eq. (9)

$$D_{m'm}^j(\pi - \phi, \pi - \theta, 0) = e^{i\pi j} (-)^{m'+m} D_{-m'm}^j(\phi, \theta, 0) \quad (13)$$

the spherical harmonics $Y_m^\ell(\theta, \phi)$ are related to the D -function via

$$D_{m0}^{\ell*}(\phi, \theta, 0) = \sqrt{\frac{4\pi}{2\ell+1}} Y_m^\ell(\theta, \phi). \quad (14)$$

The D -functions satisfy the following coupling rule:

$$D_{\mu_1 m_1}^{j_1} D_{\mu_2 m_2}^{j_2} = \sum_{j_3 \mu_3 m_3} (j_1 \mu_1 j_2 \mu_2 | j_3 \mu_3) (j_1 m_1 j_2 m_2 | j_3 m_3) D_{\mu_3 m_3}^{j_3}. \quad (15)$$

Or, equivalently,

$$D_{\mu_1 m_1}^{j_1} D_{\mu_3 m_3}^{j_3*} = \sum_{j_2 \mu_2 m_2} \left(\frac{2j_2+1}{2j_3+1} \right) (j_1 \mu_1 j_2 \mu_2 | j_3 \mu_3) (j_1 m_1 j_2 m_2 | j_3 m_3) D_{\mu_2 m_2}^{j_2*}. \quad (16)$$

Modify the formula above

$$\begin{aligned} D_{\mu_1 m_1}^{j_1} D_{\mu_3 m_3}^{j_3*} &= \sum_{j_2 \mu_2 m_2} \left(\frac{2j_2+1}{2j_3+1} \right) (j_1 \mu_1 j_2 -\mu_2 | j_3 \mu_3) (j_1 m_1 j_2 -m_2 | j_3 m_3) D_{-\mu_2 -m_2}^{j_2*} \\ &= \sum_{j_2 \mu_2 m_2} \left(\frac{2j_2+1}{2j_3+1} \right) (j_1 \mu_1 j_2 -\mu_2 | j_3 \mu_3) (j_1 m_1 j_2 -m_2 | j_3 m_3) D_{-\mu_2 -m_2}^{j_2*} \\ &= \sum_{j_2 \mu_2 m_2} \left(\frac{2j_2+1}{2j_3+1} \right) \left(\frac{2j_3+1}{2j_1+1} \right)^{1/2} \left(\frac{2j_3+1}{2j_1+1} \right)^{1/2} \\ &\quad (-)^{j_1-m_2-j_3} (j_3 \mu_3 j_2 \mu_2 | j_1 \mu_1) (-)^{j_1-\mu_2-j_3} (j_3 m_3 j_2 m_2 | j_1 m_1) (-)^{\mu_2-m_2} D_{\mu_2 m_2}^{j_2} \end{aligned} \quad (17)$$

So we have another important relationship, as an alternative to (16)

$$D_{\mu_1 m_1}^{j_1} D_{\mu_3 m_3}^{j_3*} = \sum_{j_2 \mu_2 m_2} \left(\frac{2j_2 + 1}{2j_1 + 1} \right) (j_2 \mu_3 j_2 \mu_2 | j_1 \mu_1) (j_3 m_3 j_2 m_2 | j_1 m_1) D_{\mu_2 m_2}^{j_2} \quad (18)$$

Recapitulate

$$\boxed{\begin{aligned} D_{\mu_1 m_1}^{j_1} D_{\mu_3 m_3}^{j_3*} &= \sum_{j_2 \mu_2 m_2} \left(\frac{2j_2 + 1}{2j_3 + 1} \right) (j_1 \mu_1 j_2 \mu_2 | j_3 \mu_3) (j_1 m_1 j_2 m_2 | j_3 m_3) D_{\mu_2 m_2}^{j_2*} \\ D_{\mu_1 m_1}^{j_1} D_{\mu_3 m_3}^{j_3*} &= \sum_{j_2 \mu_2 m_2} \left(\frac{2j_2 + 1}{2j_1 + 1} \right) (j_2 \mu_3 j_2 \mu_2 | j_1 \mu_1) (j_3 m_3 j_2 m_2 | j_1 m_1) D_{\mu_2 m_2}^{j_2} \end{aligned}} \quad (19)$$

It must be emphasized that the two formulas above are in fact identical. One obtains the second by taking the complex conjugate of the first and replacing the subscript 1 by 3. Using the normalization integral of (4), one obtains in addition the integrals involving three D -functions, as follows:

$$\int dR D_{\mu_1 m_1}^{j_1}(R) D_{\mu_2 m_2}^{j_2}(R) D_{\mu_3 m_3}^{j_3*}(R) = \frac{8\pi^2}{2j_3 + 1} (j_1 \mu_1 j_2 \mu_2 | j_3 \mu_3) (j_1 m_1 j_2 m_2 | j_3 m_3) . \quad (20)$$

and

$$\int dR D_{\mu_1 m_1}^{j_1}(R) D_{\mu_2 m_2}^{j_2*}(R) D_{\mu_3 m_3}^{j_3*}(R) = \frac{8\pi^2}{2j_1 + 1} (j_3 \mu_3 j_2 \mu_2 | j_1 \mu_1) (j_3 m_3 j_2 m_2 | j_1 m_1) . \quad (21)$$

Note that Eq. (21) is a complex conjugate of (20) but with indices 1 and 3 interchanged.

2 Clebsch-Gordan Coefficients

The following relationships involving Clebsch-Gordan coefficients are useful.

$$\begin{aligned}
(j_1 m_1 j_2 m_2 | j_3 m_3) &= (-)^{j_1+j_2-j_3} (j_1 - m_1 j_2 - m_2 | j_3 - m_3) \\
&= (-)^{j_1+j_2-j_3} (j_2 m_2 j_1 m_1 | j_3 m_3) \\
&= (-)^{j_1-m_1} \sqrt{\frac{2j_3+1}{2j_2+1}} (j_1 m_1 j_3 - m_3 | j_2 m_2) \\
&= (-)^{m_1-j_2+j_3} \sqrt{\frac{2j_3+1}{2j_2+1}} (j_1 - m_1 j_3 m_3 | j_2 m_2) \\
&= (-)^{m_1-j_2+j_3} \sqrt{\frac{2j_3+1}{2j_2+1}} (j_3 - m_3 j_1 m_1 | j_2 - m_2) \\
&= (-)^{j_1-m_1} \sqrt{\frac{2j_3+1}{2j_2+1}} (j_3 m_3 j_1 - m_1 | j_2 m_2) \tag{22} \\
&= (-)^{j_2+m_2} \sqrt{\frac{2j_3+1}{2j_1+1}} (j_3 - m_3 j_2 m_2 | j_1 - m_1) \\
&= (-)^{j_2+m_2} \sqrt{\frac{2j_3+1}{2j_1+1}} (j_2 - m_2 j_3 m_3 | j_1 m_1) \\
&= (-)^{j_1+m_2-j_3} \sqrt{\frac{2j_3+1}{2j_1+1}} (j_3 m_3 j_2 - m_2 | j_1 m_1) \\
&= (-)^{j_1+m_2-j_3} \sqrt{\frac{2j_3+1}{2j_1+1}} (j_2 m_2 j_3 - m_3 | j_1 - m_1)
\end{aligned}$$

The following sum rules are also useful

$$\begin{aligned}
& \sum_{m m'} (\ell' m' LM | \ell m) (\ell' m' L' M' | \ell m) \\
&= \left(\frac{2\ell + 1}{2L + 1} \right) \sum_{m m'} (-)^{\ell' - m'} (\ell m \ell' - m' | LM) (-)^{\ell' - m'} (\ell m \ell' m' | L' M') \\
&= \left(\frac{2\ell + 1}{2L + 1} \right) \delta_{LL'} \delta_{MM'} \\
& \sum_{LM} (2L + 1) (\ell' m' LM | \ell m) (\ell' n' LM | \ell n) \\
&= \sum_{LM} (2L + 1) \left(\frac{2\ell + 1}{2L + 1} \right) (-)^{\ell' - m'} (\ell m \ell' - m' | LM) (-)^{\ell' - m'} (\ell n \ell' - n' | LM) \\
&= (2\ell + 1) \delta_{m' n'} \delta_{m n}
\end{aligned} \tag{23}$$

so that

$$\begin{aligned}
& \sum_{m m'} (\ell' m' LM | \ell m) (\ell' m' L' M' | \ell m) = \left(\frac{2\ell + 1}{2L + 1} \right) \delta_{LL'} \delta_{MM'} \\
& \sum_{LM} (2L + 1) (\ell' m' LM | \ell m) (\ell' n' LM | \ell n) = (2\ell + 1) \delta_{m' n'} \delta_{m n}
\end{aligned} \tag{24}$$

The following relations involving Clebsch-Gordan coefficients can be derived by using the recursion relations for Clebsch-Gordan coefficients [Edmonds [2], p. 39]. In terms of the shorthand notations

$$\tilde{L} = L(L + 1) \quad \text{and} \quad \tilde{J} = J(J + 1)$$

one may write

$$\frac{(J - \frac{1}{2}L1|J\frac{1}{2})}{(J\frac{1}{2}L0|J\frac{1}{2})} = -\frac{2J+1}{\sqrt{\tilde{L}}} \quad (\text{odd } L \geq 1) \quad (25)$$

$$\frac{(J\frac{3}{2}L0|J\frac{3}{2})}{(J\frac{1}{2}L0|J\frac{1}{2})} = 1 - \frac{4\tilde{L}}{4\tilde{J} - 3} = (\text{even } L) \quad (26)$$

$$\frac{(J - \frac{3}{2}L2|J\frac{1}{2})}{(J - \frac{3}{2}L1|J - \frac{1}{2})} = \frac{J + \frac{1}{2}}{\sqrt{\tilde{L} - 2}} \quad (\text{even } L \geq 2) \quad (27)$$

$$\frac{(J1L0|J1)}{(J0L0|J0)} = 1 - \frac{\tilde{L}}{2\tilde{J}} = (\text{even } L) \quad (28)$$

$$\frac{(J - 1L2|J1)}{(J0L0|J0)} = -\left[\frac{\tilde{L}}{\tilde{L} - 2}\right]^{\frac{1}{2}} \quad (\text{even } L \geq 2) \quad (29)$$

$$\frac{(J2L0|J2)}{(J0L0|J0)} = 1 - \left(\frac{\tilde{L}}{2\tilde{J}}\right) \left(\frac{4\tilde{J} - \tilde{L} - 2}{\tilde{J} - 2}\right) \quad (\text{even } L) \quad (30)$$

$$\frac{(J - 2L4|J2)}{(J0L0|J0)} = \left[\frac{\tilde{L}(\tilde{L} - 6)}{(\tilde{L} - 2)(\tilde{L} - 12)}\right]^{\frac{1}{2}} \quad (\text{even } L \geq 4) \quad (31)$$

$$\frac{(J1L1|J2)}{(J0L0|J0)} = -\left[\frac{\tilde{L} - 2}{\tilde{L} - 6}\right]^{\frac{1}{2}} \left(3 - \frac{\tilde{L}}{\tilde{J}}\right) \quad (\text{even } L \geq 4). \quad (32)$$

References

- [1] M.E. Rose, Elementary theory of angular momentum (John Wiley & Sons, Inc., New York, 1957).
- [2] A.R. Edmonds, Angular momentum in quantum mechanics (Princeton University Press, Princeton, N.J., 1957).